

From Polarized Targets to Polarized Ion Beams

Opportunities and challenges for EIC spin physics



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- ① Polarized DIS with Longitudinal and Transverse nuclear polarization
 - Recent results from JLab
- ② Overview of Fixed Target Technology
- ③ Comparing polarized fixed targets with polarized ion colliders
- ④ Polarized Heavy Ions

Introduction

Polarized Deep Inelastic Scattering

$$\sigma_0 = \frac{4\alpha^2 E'^2}{q^4} \left[\frac{2}{M} F_1 \sin^2(\theta/2) + \frac{1}{\nu} F_2 \cos^2(\theta/2) \right]$$

$$2\sigma_0 A_{\parallel} = -\frac{4\alpha^2}{Q^2} \frac{E'}{E} \left[\frac{E + E' \cos \theta}{M\nu} g_1 - \frac{Q^2}{M\nu^2} g_2 \right]$$

$$2\sigma_0 A_{\perp} = -\frac{4\alpha^2}{MQ^2} \frac{E'^2}{E} \sin \theta \cos \phi \left[\frac{1}{M\nu} g_1 + \frac{2E}{M\nu^2} g_2 \right]$$

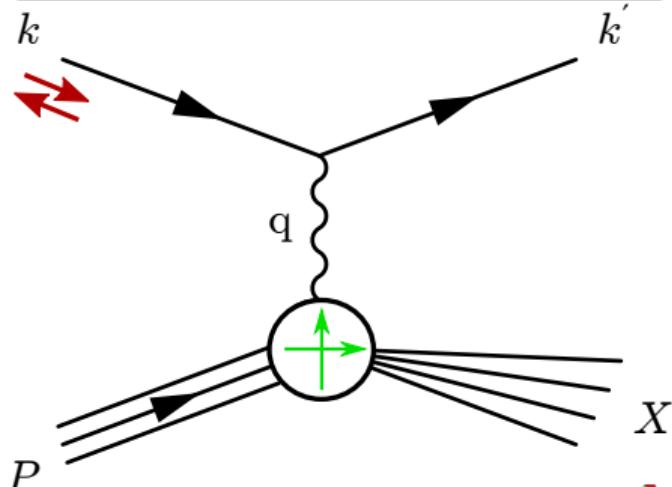
$$A_{\parallel, \perp} = \frac{A_{\parallel, \perp}^{\text{raw}}}{f P_b P_t}$$

Need \parallel and \perp polarizations

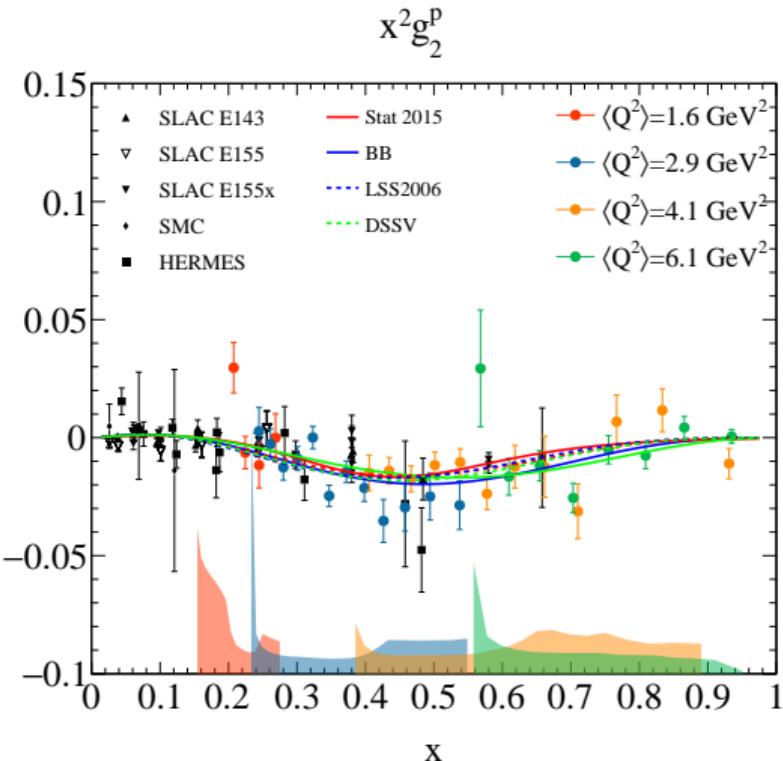
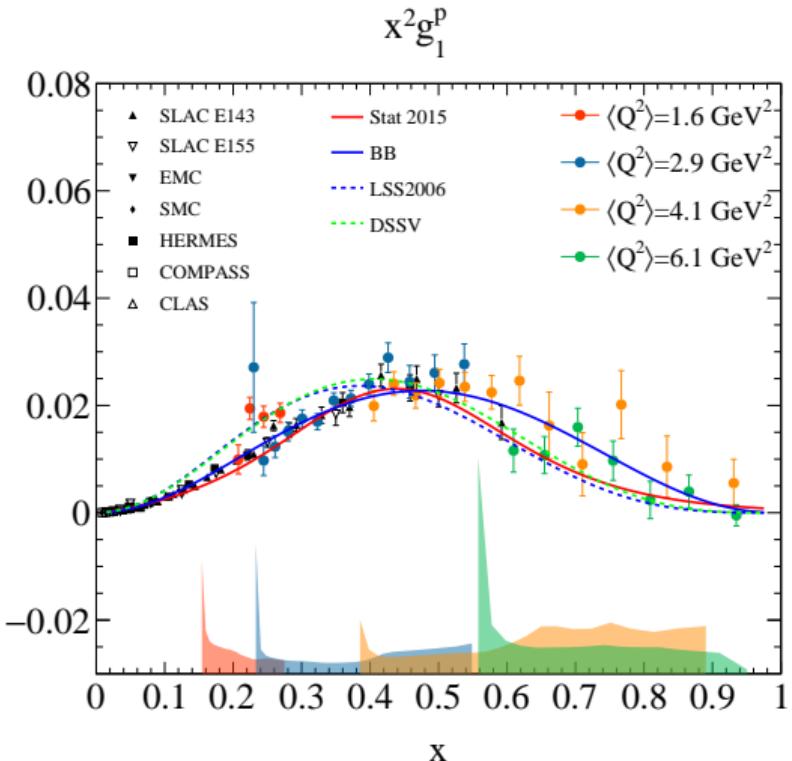
Measured Asymmetries

$$A_{\parallel}^{\text{raw}} = \frac{\sigma_{\uparrow\downarrow}^{\text{red}} - \sigma_{\uparrow\uparrow}^{\text{green}}}{\sigma_{\uparrow\downarrow}^{\text{red}} + \sigma_{\uparrow\uparrow}^{\text{green}}}$$

$$A_{\perp}^{\text{raw}} = \frac{\sigma_{\leftarrow\downarrow}^{\text{green}} - \sigma_{\leftarrow\uparrow}^{\text{red}}}{\sigma_{\leftarrow\downarrow}^{\text{green}} + \sigma_{\leftarrow\uparrow}^{\text{red}}}$$



SANE results for $x^2 g_1^p$ and $x^2 g_2^p$



The dynamical twist-3 matrix element: d_2

An average color Lorentz force

$$\int_0^1 dx x^{n-1} \left\{ g_1 + \frac{n}{n-1} g_2 \right\} = \frac{1}{2} d_{n-1} E_2^n(Q^2, g)$$

For $n = 3$

$$\int_0^1 x^2 \{2g_1 + 3g_2\} dx = d_2$$

M. Burkardt Phys.Rev.D 88,114502 (2013) and Nucl.Phys.A 735,185 (2004).

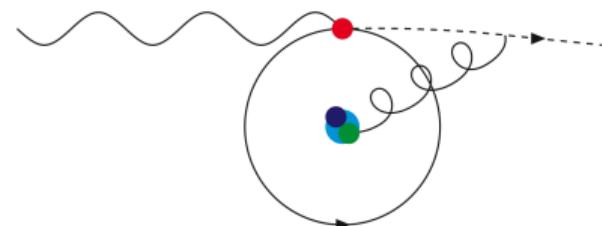
$$d_2 = \frac{1}{2MP^+{}^2 S_x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

but with $\vec{v} = -c\hat{z}$

$$\sqrt{2}G^{+y} = -E^y + B^x = -(\vec{E} + \vec{v} \times \vec{B})^y$$

Interpretations of d_2

- Color Polarizabilities (X.Ji 95, E. Stein et al. 95)
- **Average Color Lorentz force** (M.Burkardt)

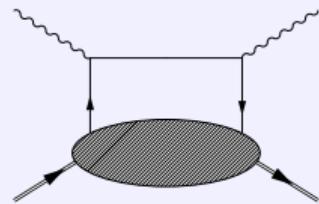


$d_2 \Rightarrow$ average color Lorentz force acting on quark moving backwards (since we are in inf. mom. frame) the instant after being struck by the virtual photon. $\langle F^y \rangle = -2M^2 d_2$

Quark-gluon Correlations : $g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$

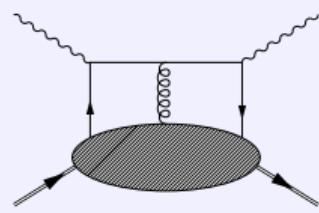
Quark-gluon Correlations : $g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$

Twist-2 (Wandzura, Wilczek, 1977)



$$g_2^{WW}(x, Q^2) = -g_1^{LT}(x, Q^2) + \int_x^1 g_1^{LT}(y, Q^2) dy/y$$
$$\equiv g_2^{tw2}(x, Q^2)$$

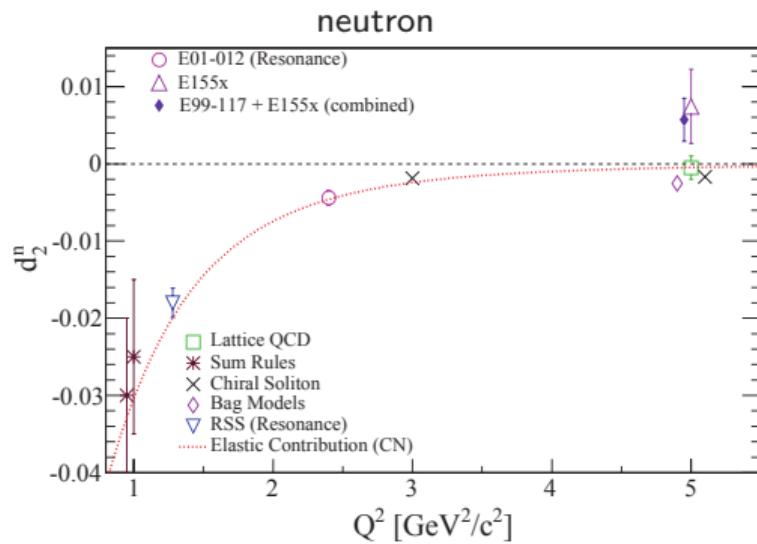
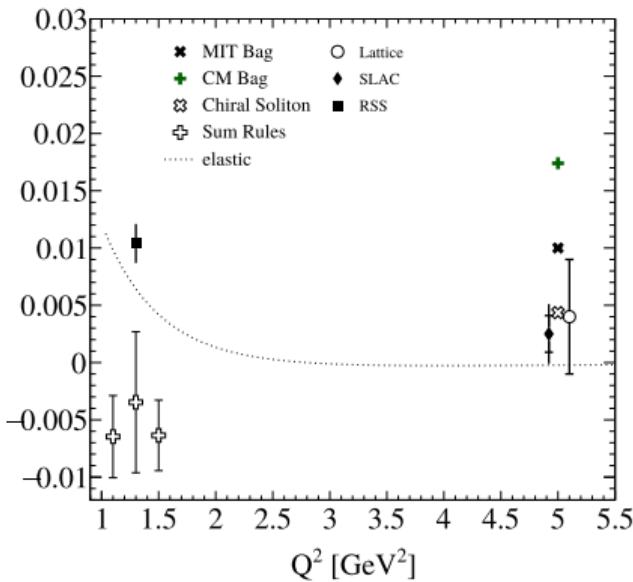
Twist-3 (Cortes, Pire, Ralston, 1992)



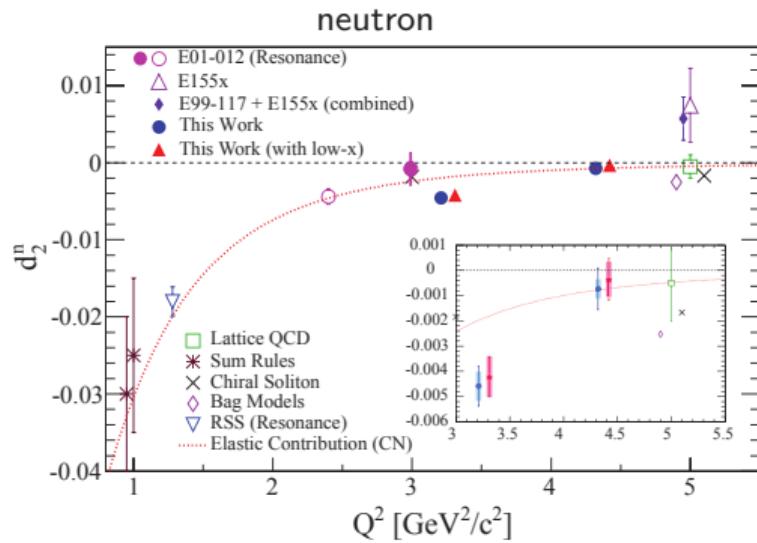
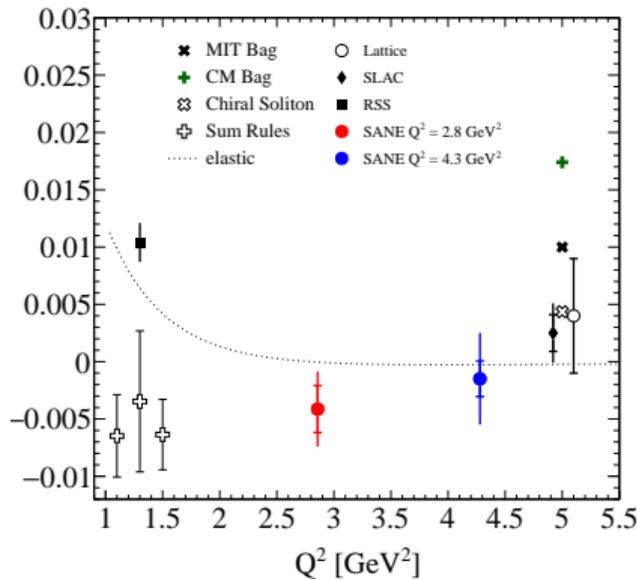
$$\bar{g}_2(x, Q^2) = - \int_{\frac{x}{M}}^1 \frac{\partial}{\partial y} \left(\frac{m_q}{M} h_T(y, Q^2) + \xi(y, Q^2) \right) \frac{dy}{y}$$
$$\equiv g_2^{tw3}(x, Q^2)$$

$$d_2(Q^2) = 3 \int_0^1 x^2 \bar{g}_2(x, Q^2) dx$$
$$= \int_0^1 x^2 (2g_1(x, Q^2) + 3g_2(x, Q^2)) dx$$

As Q^2 decreases,
when do higher twists begin to matter?
When is the color force non-zero?



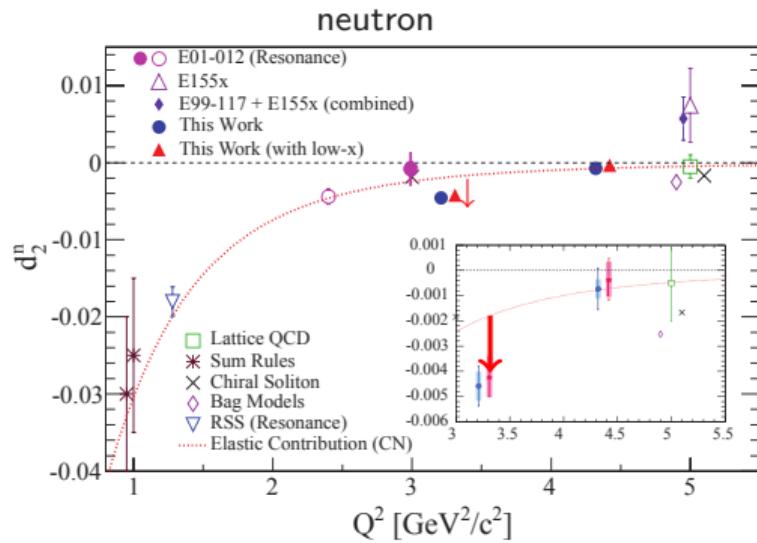
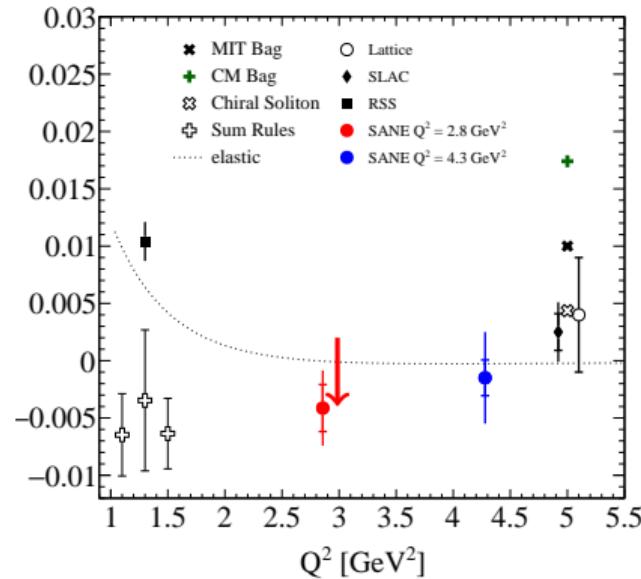
Existing data



Neutron from d_2^n experiment: D.Flay, et.al.
PRD.94(2016)no.5,052003

SANE and d_2^n Result

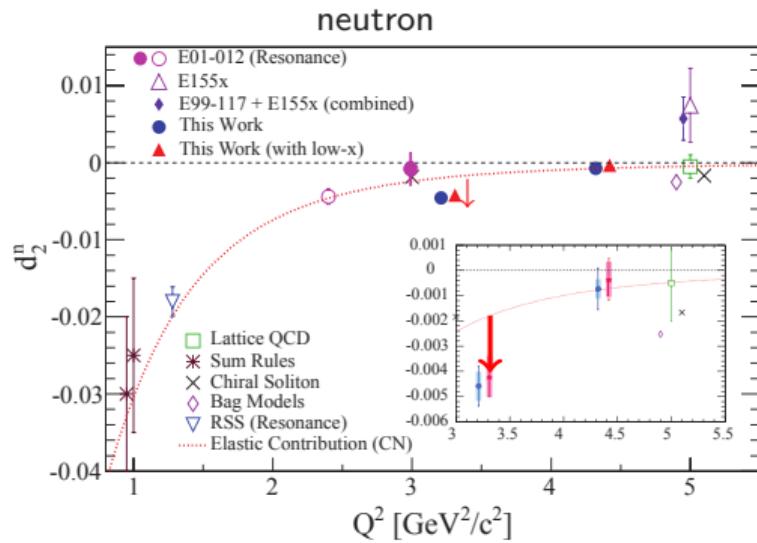
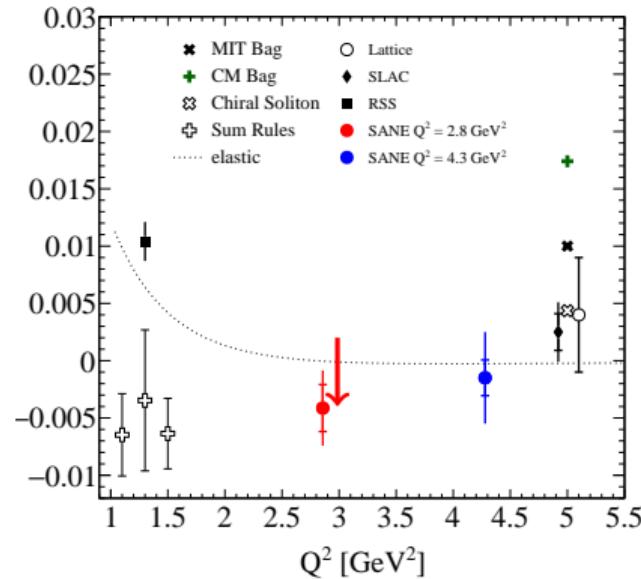
- d_2 dips around $Q^2 \sim 3$ GeV 2 for proton and neutron



Neutron from d_2^n experiment: D.Flay, et.al.
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SANE and d_2^n Result

- d_2 dips around $Q^2 \sim 3 \text{ GeV}^2$ for proton and neutron
- Is this an isospin independent average color force?



Neutron from d_2^n experiment: D.Flay, et.al.
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SANE and d_2^n Result

- d_2 dips around $Q^2 \sim 3$ GeV 2 for proton and neutron
- Is this an isospin independent average color force?
- Updated Lattice calculations are long over due!

Fixed Target Technology

A quick overview of polarized fixed targets

Dynamic Nuclear Polarization (DNP)	solid
Metastability-exchange optical pumping (MEOP)	gas
Spin exchange optical pumping (SEOP)	gas
Atomic Beam Source (ABS)	internal gas

Polarized nucleon targets

DNP \vec{p}	Solid frozen target NH ₃ , butanol, LiH
ABS \vec{p}	Internal target (Hermes)
DNP \vec{n}	From \vec{d}
SEOP \vec{n}	From ${}^3\vec{\text{He}}$
MEOP \vec{n}	From ${}^3\vec{\text{He}}$

Polarized Target Dilution Factor

Example: Polarized NH₃ Target

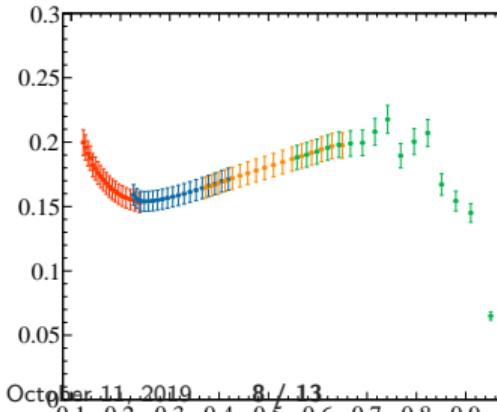
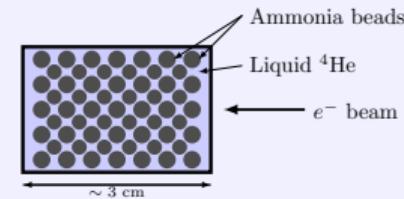
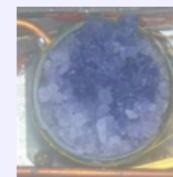
Dilution

- Takes into account scattering from unpolarized material in target.
- Need to know target geometry and material.
- Function of x and W

$$f(x, W) = \frac{N_p \sigma_p(x, W)}{N_p \sigma_p + \sum_i N_i \sigma_i(x, W)}$$

Polarized NH₃

- Packing fraction of NH₃ about 60%



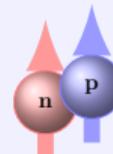
Collider Benefits

Proton



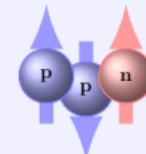
- No dilution from extra material
- No dilution from windows, cryogenics, molecular structure, ...
- Forward spectator tagging to identify struck nucleon.
- Arbitrary ion polarization direction

Deuteron



- Polarized neutron or proton

^3He



- Polarized neutron

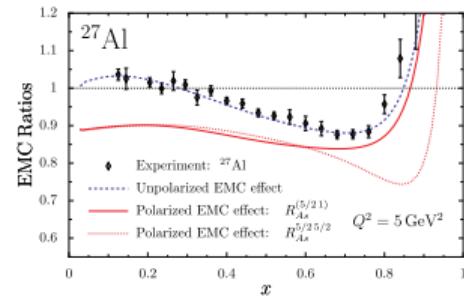
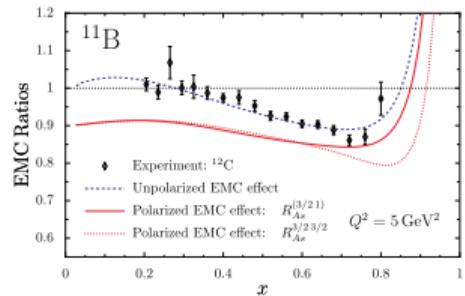
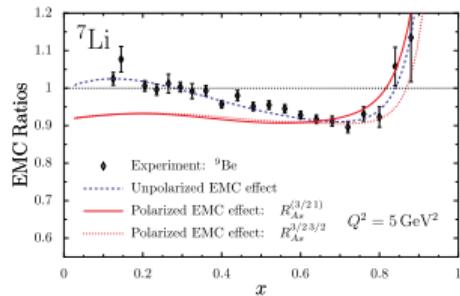
Collider Benefits

	Fixed Target	Ion Collider
Dilution	$\text{NH}_3: f \simeq 0.12$ ${}^3\text{He}: f \simeq 0.92/3$	proton: no dilution neutron: $f \simeq 1/3$
Spectator Tagging	Very difficult	Possible with forward detectors
Luminosity	$\text{NH}_3:$ Beam current limited to 100 nA $\rightarrow \mathcal{L} \simeq 10^{35} \text{s}^{-1} \text{cm}^{-2}$ ${}^3\text{He}:$ $\mathcal{L} \simeq 10^{37} \text{s}^{-1} \text{cm}^{-2}$	$\mathcal{L} \simeq 10^{34} \text{s}^{-1} \text{cm}^{-2}$ better dilution compensates for lower luminosity
\parallel, \perp polarization	$\text{NH}^3:$ physically rotated 5T magnet leads to different rates/backgrounds in detectors for same kinematics ${}^3\text{He}:$ weak field, dual Helmholtz coils for easy rotation.	Bunch by bunch ion spin rotation?

Polarized Heavy Ions

Polarized EMC Effect

$$R \simeq g_1^A / g_1^p$$



Cloët, et.al., Phys.Rev.Lett. 95 (2005) 052302

Tagging to identify struck system

- Full tagging of spectator system (A-1)
- Identify struck nucleon to eliminate dilution of nucleus
- Would like many polarized ions beyond ${}^3\text{He}$

Laser Driven Source

25 Years Ago at Argonne

PROSPECT OF POLARIZED TARGETS IN ELECTRON RINGS*

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The feasibility of performing experiments with polarized targets in electron rings is discussed. Examples of the physics which would be accessible by this method are given. It is noted that this method is consistent with recent proposals for linear colliders. A new method for producing a polarized atomic beam is proposed and some preliminary calculations are presented. A summary of laser-driven polarized atomic beams is also given.

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Nuclear Instruments and Methods in Physics Research A 364 (1995) 58–69

NUCLEAR INSTRUMENTS & METHODS IN PHYSICS RESEARCH
Section A

Laser-driven source of spin-polarized atomic hydrogen and deuterium

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Received 20 January 1995

An active storage cell for a polarized gas internal target

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Received 8 June 1994

Recent Developments

- Hybrid SEOP → K and Rb (M.V. Romalis PRL 105, 243001 (2010))
- Readily available high power diode lasers for pumping Rb (795 nm)
- Successful polarized ^3He program at JLab.

Beginning to investigate general purpose hybrid SEOP to polarize heavier ions such as ^{21}Ne .

Summary

- Nuclear polarization is key for unraveling QCD at the EIC
- **All polarization directions equally important**, especially for imaging program
- Extreme forward tagging will **significantly improve the science** extracted with each polarized ion electron collision
- Nuclear polarization is needed to investigate Polarized EMC Effect
- A general purpose laser driven source may provide polarized heavy ions

Thank You!

Backup

E07-003 : Big Electron Telescope Array

